

Title	Mechanism of thermal atomic layer etch of w metal using sequential oxidation and chlorination: A first-principles study
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Publication date	2020-07-15
Original Citation	Kondati Natarajan, S., Nolan, M., Theofanis, P., Mokhtarzadeh, C. and Clendenning, S. B. (2020) 'Mechanism of Thermal Atomic Layer Etch of W Metal Using Sequential Oxidation and Chlorination: A First-Principles Study', ACS Applied Materials & Interfaces, 12(32), pp. 36670-36680. doi: 10.1021/acsami.0c06628
Type of publication	Article (peer-reviewed)
Link to publisher's version	https://pubs.acs.org/doi/10.1021/acsami.0c06628 - 10.1021/acsami.0c06628
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Mechanism of Thermal Atomic Layer Etch of W Metal Using Sequential Oxidation and Chlorination: A First-Principles Study

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Cite This: *ACS Appl. Mater. Interfaces* 2020, 12, 36670–36680



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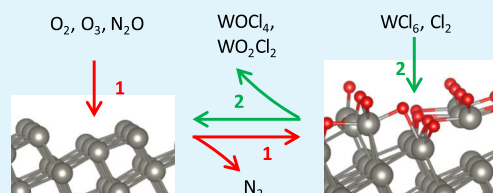
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Supporting Information

ABSTRACT: Thermal atomic layer etch (ALE) of W metal can be achieved by sequential self-limiting oxidation and chlorination reactions at elevated temperatures. In this paper, we analyze the reaction mechanisms of W ALE using the first-principles simulation. We show that oxidizing agents such as O_2 , O_3 , and N_2O can be used to produce a WO_x surface layer in the first step of an ALE process with ozone being the most reactive. While the oxidation pulse on clean W is very exergonic, our study suggests that runaway oxidation of W is not thermodynamically favorable. In the second ALE pulse, WCl_6 and Cl_2 remove the oxidized surface W atoms by the formation of volatile tungsten oxychloride ($W_xO_yCl_z$) species. In this pulse, each adsorbed WCl_6 molecule was found to remove one surface W atom with a moderate energy cost. Our calculations further show that the desorption of the additional etch products is endothermic by up to 4.7 eV. Our findings are consistent with the high temperatures needed to produce ALE in experiments. In total, our quantum chemical calculations have identified the lowest energy pathways for ALE of tungsten metal along with the most likely etch products, and these findings may help guide the development of improved etch reagents.

KEYWORDS: atomic layer etching, transistor contacts, density functional theory, self-limiting reaction, atomistic simulations, first principles



1. INTRODUCTION

Atomic layer etch (ALE) processing has gained considerable attention in recent years within the semiconductor industry.^{1–4} The continued scaling of semiconductor devices demands the use of ever-thinner and higher-performing materials, which drives the development of gentle etch process technologies with atomic-level precision and high degrees of selectivity toward surrounding materials. ALE shares many similarities with the well-known atomic layer deposition (ALD), widely used in the semiconductor industry, and can be considered as the reverse of ALD. ALE permits the removal of thin films layer by layer exploiting atomic-scale precision inherent in the method using sequential and self-limiting surface reactions,^{1–3,5–7} similar to ALD. Currently used ALE processes are anisotropic using directional high energy ion bombardment to drive the removal of the modified surface layers.^{2,8} On the other hand, a thermal ALE process relies on temperature and thermochemically favorable reactions to drive the removal of surface species.⁹ There are a multitude of examples for thermal ALE of various materials, including oxides (HfO_2 ,^{3,7,10,11} ZrO_2 ,^{3,11} SiO_2 ,¹² WO_3 ,¹³ ZnO ,¹⁴ and Al_2O_3 ,^{9,11,15–18}), nitrides (including AlN ,¹⁹ GaN ,²⁰ and TiN ,^{21,22}), fluorides (e.g., AlF_3 ,²³), and some work on W.^{13,24–26} Other ALE techniques include plasma ALE of SiO_2 ,^{27,28} ZnO ,²⁹ and GaN ,^{30,31} and infrared annealing ALE of Si_3N_4 .³² Despite this effort, the

atomistic details of the mechanism of thermal ALE processes are still not fully understood.

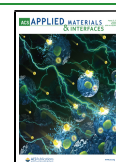
Tungsten, because of its low resistivity, resistance to electromigration, and ease of thin film growth by atomic layer and chemical vapor deposition, is often used in source/drain contacts in CMOS devices.³³ Unlike ALD processes for W, which have received considerable attention,^{34–42} the utility of a complementary W atomic layer etch process has only recently emerged.^{13,24–26,43} Due to the continued scaling of critical dimensions in next-generation semiconductor devices and the incorporation of a more varied material palette in accordance with Moore's law, the need for novel digital etch methods to enable the precise removal of thin film materials is of extreme importance.^{2,44}

While there are anisotropic W ALE processing methodologies available,⁴⁵ isotropic thermal ALE techniques are required to perform conformal etch in high aspect ratio structures. Indeed, progress toward the isotropic thermal ALE of metallic W has been reported by the groups of George and

Received: April 10, 2020

Accepted: July 15, 2020

Published: July 15, 2020



Parsons. The work reported by George *et al.*¹³ has shown the quasi-ALE of metallic W via a conversion etch mechanism utilizing ozone (O₃), BCl₃, and anhydrous HF vapor. In this system, ozone was used to oxidize the W surface to WO₃ in a diffusion-limited process, which then reacts with BCl₃ to generate a volatile etch product in the form of WO_xCl_y, while tandemly generating a B₂O₃ surface layer, which is then susceptible to self-limiting etch by HF. Parsons and co-workers²⁴ have recently reported the ALE process of W with sequential use of O₂ plasma or O₃ and WF₆. The oxidizing plasma produces a diffusion-limited WO₃ surface layer, which can be removed by exposure to WF₆ yielding WO₂F₂ as a volatile etch byproduct. They have also reported results on an oxidation plus WCl₆-based W ALE process^{25,26} in which they used a thermodynamic analysis to rationalize the experimental data. Despite these recent advances in atomic layer etch, however, mechanistic insight into the ALE processes remains limited. Computational investigations of a thermal ALE process for metal oxides with Al₂O₃ as an example have been reported previously,⁴⁶ but a similar investigation for metal ALE is lacking in the literature. This is due to an abundance of reactive transient species and etch products in a given ALE cycle and the difficulties related to *in situ* metrologies needed to identify and quantify these intermediates and products. Accordingly, to enable the rapid development of such processes, it is imperative that a mechanistic understanding of W ALE be established.

Herein, we report a first-principles computational study using density functional theory (DFT) to explore an oxidation/chlorination thermal ALE mechanism for W metal. In this study, the nature of the surface oxidation step to produce a WO_x layer on tungsten is investigated using O₂, O₃, and N₂O as oxidants to uncover potential routes toward self-limiting oxidation of the W surface. Second, WCl₆ and Cl₂ were examined as potential etch reagents to generate volatile WO_xCl_y etch species. Spin-polarized DFT coupled with density functional perturbation theory (DFPT) for a thermochemical analysis⁴⁷ were used to address self-limiting versus continuous etching pathways. At the same time, ground state spin polarized DFT was used to calculate the energies of surface species along several proposed reaction pathways. From these results, a full atomic layer etch chemical pathway for W metal can be mapped yielding a feasible reaction mechanism featuring sequential oxidation and chlorination steps.

2. COMPUTATIONAL SETUP

All calculations reported in this paper are based on spin-polarized DFT using VASP v5.4.⁴⁸ Core electrons are represented by projector augmented wave (PAW) potentials,^{49,50} and valence electrons are treated explicitly by expanding their wave functions in a plane wave basis set with an energy cutoff of 400 eV. The exchange and correlation (XC) contributions are approximated by the generalized gradient approximation (GGA) using the Perdew–Burke–Ernzerhof (PBE) functional.⁵¹ The total energies and forces of the geometries are converged within 1×10^{-4} eV and -0.02 eV/Å, respectively. Methfessel–Paxton first-order smearing with 0.1 eV broadening is used for the scf calculations. The effect of dispersion correction on the structure and energies of bulk and surface geometries was found to be insignificant for this system, and thus dispersion corrections are not included (see section S0 of the Supporting Information).

The reported reaction free energies are computed as follows:

$$\Delta G = \Delta H - T\Delta S + RT\ln(Q) \quad (1)$$

$$\Delta H = \Delta E + \Delta ZPE + \Delta W(T) \quad (2)$$

$$Q = \prod p_{\text{products}}^{\mu} / \prod p_{\text{reactants}}^{\mu} \quad (3)$$

Here, ΔG , ΔH , and ΔS correspond to reaction free energy, reaction enthalpy, and reaction entropy, respectively. ΔE , ΔZPE , and $\Delta W(T)$ refer to the electronic reaction energy at 0 K, zero point energy change, and temperature-dependent enthalpy change, respectively. Q is the reaction quotient used to include reactant and product pressures. A reactant pressure of 0.2 Torr and a product pressure of 0.01 Torr are used for the reported free energy calculations (see section S7 of the Supporting Information for other pressure combinations). Pressure contributions to the free energy are included only for the gas phase species. R and μ are the gas constant and stoichiometric coefficient, respectively. The H and S values of the bulk and surface geometries are computed using the Phonopy code,⁵² which requires accurate force constants obtained from DFPT calculations in VASP with a strict energy convergence threshold of 1.0×10^{-8} eV. However, for convenience, the H and S values for gas phase molecules are obtained from the Turbomole package version 6.2⁵³ using the PBE XC functional and the def-TZVPP basis set.

W has a BCC crystal structure⁵⁴ with a calculated lattice constant of 3.17 Å, which agrees with the experimental value of 3.16 Å.⁵⁵ The cohesive energy of W is 9.1 eV/atom, which compares well to the experimental value of 8.9 eV/atom.⁵⁶ For the bulk W calculation, a $12 \times 12 \times 12$ Monkhorst-Pack K-point sampling mesh is used and the lattice constant is obtained by simultaneously relaxing the ionic positions, cell shape, and cell volume at a larger plane wave energy cutoff of 550 eV. For the surface calculations, we have chosen the (2×3) supercell of a 10.8 Å thick stepped W(3 1 0) surface that has a computed surface energy of 3.77 J/m². The high index high energy W stepped surface was chosen so as to maximize reactivity to incoming reagents as compared to the flat, less reactive low index surfaces. This metal slab consists of six W bilayers with 12 atoms each (72 W atoms in total) with a surface area of 0.955 nm². A vacuum of 19 Å separates the two surfaces of the slab along the surface normal direction. A $3 \times 3 \times 1$ Monkhorst-Pack K-point sampling mesh is used for all surface calculations to account for the supercell size used in this study. Atomistic models of bulk W and W(3 1 0) surface are shown in Figure S1 of the Supporting Information.

Tungsten sites are 8-fold coordinated in the bulk due to the BCC crystal structure. Each bulk W atom has six valence electrons shared between eight neighbors; therefore, three-fourths of an electron is contributed to each metallic W–W bond. In the (2×3) supercell of W(3 1 0), there are 12 surface W atoms of which 6 are 4-fold coordinated and the other 6 are 6-fold coordinated. This amounts to 27 unpaired electrons on the surface. Therefore, we adsorbed 14 O atoms (coverage of 14.7 O/nm²) per supercell to the bare surface to represent an oxidized W surface region, which is computationally tractable. The degree of oxidation will depend on processing conditions (Section 3.2), but this provides a reasonable model of partially oxidized W for our analysis. We considered different arrangements of O species on the surface and relaxed the geometries (see Figure S2 of the Supporting Information). The most stable WO_x structure we have found is shown in Figure 1a,b. This geometry represents

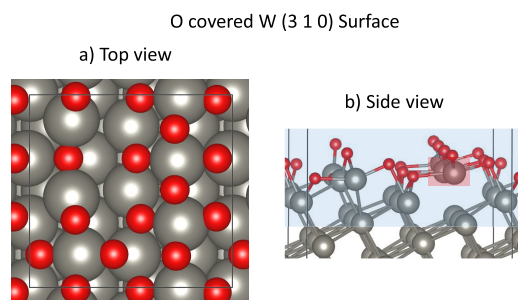


Figure 1. (a) Top and (b) side views of the most stable O-covered W(3 1 0) surface. The blue region indicates the atoms used in the DFPT calculation. The top W atoms coordinated by more than two O atoms, highlighted in the red region, are lifted out of their lattice positions due to the formation of W–O bonds. The color coding is as follows: W = gray, O = red.

the partially oxidized W surface that is present after the first ALE (oxidation) pulse. From Figure 1a,b, we find that the O atoms prefer to adsorb at the bridge and 3-fold sites on the surface. Moreover, the surface W–O bonds would weaken the underlying W–W bonds, which is evident from the out-of-lattice positions of the surface W atoms highlighted in red.

3. RESULTS

3.1. Volatile Products of Interest. The proposed W-etch species, which originate from the surface for this ALE process, have been determined to be either WOCl_4 or WO_2Cl_2 ^{24–26} depending on the etch gas stoichiometry. To etch a single atom of W from the surface (W_s) with a single molecule of WCl_6 , the incoming WCl_6 molecule must react with the surface to generate a minimum of two volatile W species in the form of WOCl_4 or WO_2Cl_2 . In order to maintain proper “book-keeping” of volatile species, it is important to note that the first volatile W species corresponds to the W atom originating from the incoming WCl_6 molecule. Due to the origin of this W atom from the reactant gas, the resulting volatile W species (either WOCl_4 or WO_2Cl_2) is considered a non- W_s etch species (NES), while the second W-containing species would be designated the first- W_s etch species (FES). In the most favorable case, one incoming WCl_6 could react with the surface to form 3 equiv of WO_2Cl_2 , thereby effectively etching two surface W atoms. The first molecule of WO_2Cl_2 would be a NES, while the first and second W species that originate from the surface would be considered the first- W_s etch species (FES) and second- W_s etch species (SES), respectively. This nomenclature is used throughout the remainder of this paper. Alternatively, a second mechanism utilizing 2 equiv of WCl_6 to remove one surface W atom resulting in the formation three WOCl_4 species is also considered. The first two W species, which originate from the incoming WCl_6 molecule, would be considered NESs, while the last W species, which originates from the surface, would be a FES.

If we consider Cl_2 as an etch gas, each molecule of Cl_2 has the potential to remove an oxidized surface W atom in the form of WO_2Cl_2 (a FES). Likewise, 2 equiv of Cl_2 could be used to form 1 equiv of WOCl_4 originated from the surface (a FES). A schematic showing the described volatile byproducts is shown in Figure S3 of the Supporting Information. It is important to note that the formation energies for these proposed volatile species were very similar (WOCl_4 : -7.1 eV; WO_2Cl_2 : -7.3 eV) using bulk W, O_2 , and Cl_2 molecules as a

reference. However, their formation and desorption energies with respect to a reactive surface differ and will be addressed later in this text. We first study how the oxidizing agents interact with the bare surface to form the WO_x surface shown in Figure 1a and then examine how WCl_6 and Cl_2 molecules interact with this oxidized W surface.

3.2. Mechanisms in the Oxidation Pulse. In the first ALE pulse, we modify the W surface by oxidizing it with O_2 , O_3 , and N_2O . N_2O interaction with the bare W surface is the weakest when compared to O_2 and O_3 ; therefore, the discussion on N_2O –W interaction is moved to section S2 of the Supporting Information. Oxidation of W using O_2 and O_3 has been studied quite extensively, and it is known that bulk oxidation of W in an O_2 atmosphere is appreciable only above 500°C , which is well above the ALE processing temperature.^{57–59} An account on the temperature dependence of the oxidation of W using O_2 and O_3 relevant to atomic layer etching has been discussed by George and co-workers.¹³ From the N-E analysis⁴⁷ given in section S3 of the Supporting Information, comparing the continuous and self-limiting reaction models of the first ALE pulse, we predicted a minimum (thermodynamic) energy barrier of about 2.3 eV to cause bulk oxidation. We could expect that the predicted thermodynamic barrier to be breached below 800°C as W oxide, formed on the surface, is reported to become volatile beyond that temperature.⁶⁰ Warren *et al.*⁶¹ found that a 1 h exposure of polycrystalline W to O_2 below 200°C resulted in the formation of 1.0–1.6 nm of WO_3 . Therefore, by keeping the temperature below 200°C or by controlling the exposure time and partial pressure of the reactants, a thin layer of oxide relevant for ALE can be formed. The temperature needed for the thermal ALE process will also be influenced by the activation energies required in the second ALE pulse.

On the clean W(3 1 0) surface, both O_2 and O_3 molecules dissociate spontaneously on adsorption and form surface W–O bonds as shown in Figure 1a,b by releasing -7.24 and -12.13 eV per adsorbed molecule, respectively. For comparison, the most favorable binding energy of a N_2O molecule is -5.81 eV. O_3 is a stronger oxidant than the O_2 molecule. Since the adsorption geometries and binding sites are similar for both molecules, we will only examine the adsorption of the O_2 molecule in detail. We considered different binding sites for the dissociated O atoms on the W surface as shown in Figure 2a–e. The binding energies and shortest W–O bond distances from these geometries are listed in Table 1. The computed binding energy per O atom is in the range of -3.30 to -4.47 eV for on-surface adsorption, with the bridge sites being the most favorable followed by the atop, 3-fold, and 4-fold sites. The adsorption at subsurface sites is about 3 eV less favorable than the most favorable bridge site (Figure 2f). Clearly, at this coverage of $2\text{ O}/\text{nm}^2$, diffusion to subsurface sites is not as favorable as on-surface adsorption. The O atom in the subsurface site is at least 1.9 \AA away from the neighboring W atoms. The shortest W–O bond length is found at the atop site followed by the bridge, subsurface, 3-fold, and 4-fold sites. Comparing with the most favorable oxygen-covered surface model shown in Figure 1a, we note that the O atoms are predominantly adsorbed at the bridge sites.

In this section, we saw the various adsorption sites preferred by the dissociated O atoms and that the O_2 molecule indeed dissociates spontaneously. We have prepared the saturated surface models as described in the computational section. Any additional O_2 molecule introduced in the system beyond a

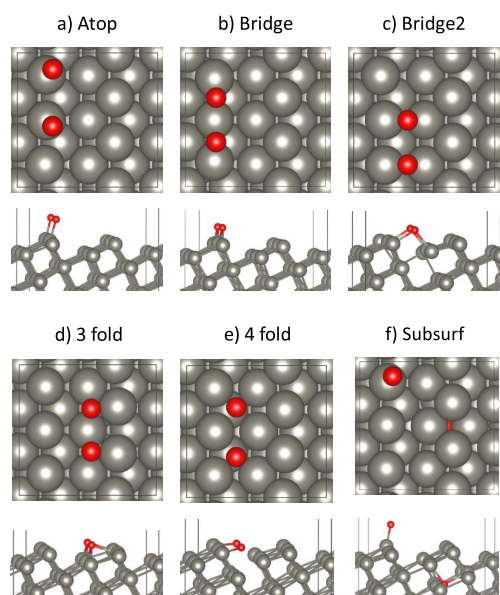


Figure 2. The top and side views of dissociated O_2 molecules on the $\text{W}(3\ 1\ 0)$ surface at atop sites, bridge sites, alternate bridge sites, 3-fold sites, 4-fold sites, and subsurface sites are shown in (a), (b), (c), (d), (e), and (f), respectively. The color coding is as follows: W = gray, O = red.

Table 1. Binding Energies and Minimum W–O Bond Distance for O_2 adsorption on the $\text{W}(3\ 1\ 0)$ Surface^a

geometry	E_{bind} (eV/O)	$d(\text{W}-\text{O})$ (Å)
atop	−3.80	1.7
bridge	−4.47	1.9
bridge2	−3.99	1.9
3-fold	−3.31	2.0
4-fold	−3.30	2.1
subsurface	−1.45	1.7 (1.9)

^aFor the subsurface geometry, the distance in parentheses correspond to the O atom in the subsurface site.

single O_2 molecule described in this section will indeed dissociate (in an exoergic reaction) until the saturated surface is formed. Thus, the intermediate steps do not need to be discussed.

3.3. Mechanisms in the Chlorination Pulse. The primary mechanism of chlorination of the oxidized W surface by WCl_6 is the dissociation of the molecule and subsequent exchange of Cl ligands from the reagent to the substrate, which requires elevated temperatures. These are complex dynamical processes unlike the simple dissociative adsorption of the reactant molecules in the first pulse. Before investigating the chlorination of the oxidized W surface, we first explored the possibility of uncontrolled etching of W metal by WCl_6 and Cl_2 . These results are discussed in detail in section S4 of the Supporting Information. The exposure of WCl_6 and Cl_2 on the bare W surface did not result in continuous etching of the substrate even at elevated temperatures. Deposition of W, when WCl_6 is used, is also not thermodynamically favorable beyond room temperature.

WCl_6 being an octahedral molecule interacts weakly with the O-covered W surface due to the net electrostatic repulsion between the O and Cl species. As a result of this net electrostatic repulsion, WCl_6 does not dissociate spontaneously after initial adsorption at low temperatures due to the presence

of the kinetic energy barrier. Dissociation of the WCl_6 molecule would need elevated temperatures to excite the W–Cl bonds or O-free regions on the surface as WCl_6 readily dissociates on the bare W surface. Experiments suggest that a temperature of at least 200 °C is needed for the ALE process to take place,^{25,26} which means that the WCl_6 molecules would have dissociated at this temperature. A short 2.5 ps duration ab initio molecular dynamics (MD) simulation at 800 K of the WO_x surface with two physisorbed WCl_6 molecules revealed a spontaneous donation of one of the Cl ligands by each of the reactant molecules to the surface W atoms as shown in Figure 3. A high temperature of 800 K is used to accelerate the MD simulation; otherwise, we would have needed a significantly longer simulation time to observe this reaction step at lower temperatures.

The W atom in the WCl_6 molecule is the donor W atom (W_d), while the acceptor W (W_a) is one of the oxidized surface W atoms. In the first 300 steps (750 fs) of the simulation, we can identify the high amplitude vibrational mode of the Cl atom from the fluctuations in the interatomic W_d –Cl distances (see Figure 3). Moreover, an increase in the W_d –Cl distance correlates with a decrease in the W_a –Cl distance. This large amplitude motion, a byproduct of the thermal energy contribution, results in the dissociation of one of the Cl ligands from the WCl_6 molecule. After the crossover point at 837 fs (step 335), the dissociated Cl ligand did not attempt to bind back to the W_d atom, which is evident from the continuous increase in $d(\text{W}_d$ –Cl) value beyond 1.25 ps (step 500). The inset graph in Figure 3 shows the distances and single-point energies of the system during the dissociation of WCl_6 . We see an energy change of 800 meV right before the dissociation takes place; this includes contribution from the motion of surface atoms as well. The dissociation of the Cl ligand from the second WCl_6 molecule was similar to the above observation.

To compute the entire dissociation pathway, the formation of non-volatile species are considered when WCl_6 is adsorbed either as an intact or dissociated molecule, for which the geometries are shown in Figure 4A–E. While there are very many possibilities to arrange an intact and dissociated WCl_6 molecule on the O-covered W surface, it is not practical to study them all; therefore, we have considered these five minimum geometries as a representative for further discussion. In Figure 4, geometry A shows a weak interaction between the WCl_6 molecule and the O-covered W surface with a binding energy of just −0.09 eV. Geometry B shows the state where the weakly adsorbed WCl_6 dissociates and a Cl ligand is adsorbed to a surface W atom. This reaction is not spontaneous and costs 0.48 eV. The potential energy surface cut along the pathway, obtained by a linear interpolation of the position of the dissociated Cl ligand from geometry A to B, is computed, and the computed dissociation barrier is 0.52 eV (graph and geometries in Figure S9 of the Supporting Information). The reverse barrier for the reformation of WCl_6 is just 0.04 eV. However, this reaction is hindered by the diffusion of the WCl_5 fragment into the vacuum immediately after the dissociation. It is possible that all impinged WCl_6 molecules follow this process and chlorinates the WO_x surface without depositing any W as WCl_5 . However, the WCl_5 fragment in the gas phase can further adsorb on the surface by forming W–O bonds as shown in geometry C (WCl_5 adsorbed), which is energetically favorable and releases −1.32 eV of energy in comparison to geometry B. Dissociation of a Cl ligand from the adsorbed

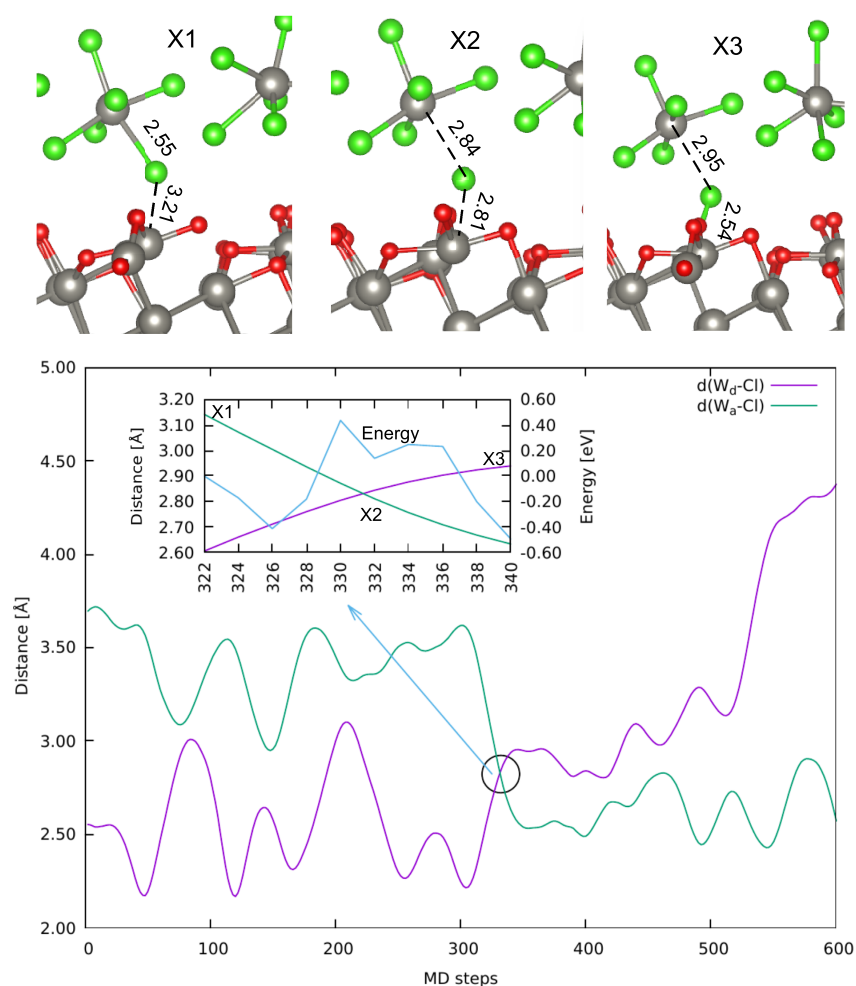


Figure 3. Snapshots (X1–X3) from MD simulation of two WCl_6 molecules adsorbed on the WO_x surface along with the evolution of distances between the dissociated Cl atom and the donor/acceptor W atoms. The inset graph shows the distances and single-point DFT energies of the snapshots (relative to that of snapshot 322, X1) during the bond dissociation. The color coding is as follows: W = gray, O = red, and Cl = green.

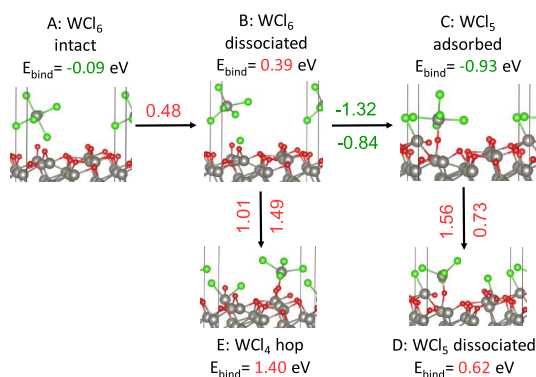


Figure 4. Relaxed geometries showing the non-volatile (dissociative) adsorption of WCl_6 on the O-covered W surface. The values in red indicate an energy cost, and the values in green indicate a gain in energy at 0 K. The values above (or left of) the arrow are the energy change with respect to the preceding configuration, while the values below (or right of) the arrow are the energy change with respect to geometry A (cumulative energy change). In geometry A, WCl_6 is intact, and in geometries B, C, D, and E, the molecule is in the dissociated state. The reference gas phase molecule used for the binding energy calculation is WCl_6 .

WCl_5 fragment in geometry C to form geometry D has an energy cost of 1.56 eV, suggesting that further dissociation of

Cl ligands from surface-adsorbed WCl_5 has a large energy penalty. Dissociation of surface-bound WCl_5 is unfavorable at 0 K due to enthalpic loss, which will be compensated by entropic gain at elevated temperatures due to the formation of the volatile species (WOCl_4). Another possibility for the WCl_5 fragment in geometry B is to form geometry E by donating another Cl ligand in the gas phase to the surface and adsorb as WCl_4 by forming a W–O bond (WCl_4 hop). This results in the formation of WOCl_4 , one of the proposed volatile species, which spontaneously desorbs from the surface, as shown in geometry E. Once again, this step will see an entropic gain at elevated temperatures and compensate the enthalpic loss reported at 0 K. Only geometries A and C have favorable (negative) binding energies at 0 K.

Beginning from the most favorable non-volatile adsorption geometry C, step-by-step desorption of the proposed volatile etch products and their corresponding energetics are investigated as shown in Figure 5. Three cases in which the sequential removal of 1 WOCl_4 + 1 WO_2Cl_2 , 1 WO_2Cl_2 + 1 WOCl_4 , and 3 WO_2Cl_2 are considered. Note that the first volatile product removed in all three cases is the NES followed by FES and SES as discussed earlier. In the first case, WOCl_4 was removed as the NES, with a desorption energy of 0.44 eV. Due to entropic gain, this step is favorable (−1.73 eV gain) at 500 K. This however does not count toward etching of the

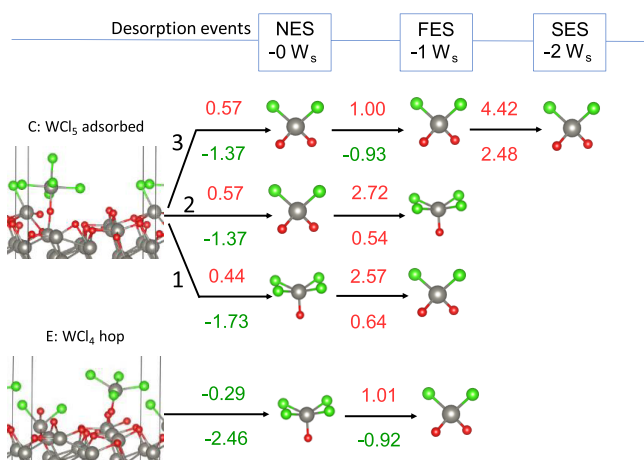


Figure 5. Energy requirements for the removal of volatile species from geometry C and E from Figure 4. The values in red indicate a moderate energy cost, and the values in green indicate a gain in energy. The values above the arrow are the desorption energies of the respective volatile species, while the values below are the free energy change of desorption at 500 K. The desorbed molecules are either WOCl₄ or WO₂Cl₂.

substrate since the W atom removed is not a surface atom. Following this step, WO₂Cl₂ is removed as the FES with an additional desorption energy of 2.57 eV (which is also not favorable at 500 K). In the second case, WO₂Cl₂ is removed first as the NES with an energy cost of 0.57 eV, which is 0.13 eV more than WOCl₄ in the first pathway. This is because a second W–O bond must be broken in forming WO₂Cl₂. The subsequent removal of WOCl₄ in this sequence costs 2.72 eV of energy. We note that WOCl₄ offers more entropy than WO₂Cl₂ due to the presence of an additional ligand, and it is evident from the reduced free energy gain in the first step (−1.37 eV for WO₂Cl₂ against −1.73 eV for WOCl₄) and reduced free energy loss in the second step (0.54 eV for WOCl₄ against 0.64 eV for WO₂Cl₂) when compared with the previous case. In these two cases, the removal of the FES is not favorable even at 500 K. However, the cumulative enthalpic gain in the earlier steps of the ALE cycle could offset this energy cost.

In the third case, three WO₂Cl₂ molecules were removed sequentially. While the removal of the NES costs 0.57 eV, the removal of the FES and SES becomes increasingly expensive with energy costs of 1.00 and 4.42 eV, respectively. In contrast to the first two cases, the removal of the FES in this case is favorable at 500 K (−0.93 eV); however, the removal of the SES is not. This is due to the fact that a considerable number of surface W–O and W–W bonds have to be broken to remove the second WO₂Cl₂ species. Similarly, from the minimum geometry E, the removal of the pre-formed non-W_s etch product WOCl₄ is favorable by −0.29 eV. Further removal of the etch product WO₂Cl₂ costs an energy of 1.0 eV, but this step is entropically favorable at 500 K in contrast to the first case from geometry C. The free energy values for other reactant and product pressures for comparison is given in section S7 of the Supporting Information. It is to be noted that the influence of the reactant pressure to the free energy is very small compared to the contributions from enthalpy and entropy in this case.

From the above analysis, we find that each WCl₆ molecule is capable of removing up to two surface W atoms and six surface

O atoms in the best case scenario (3 WO₂Cl₂). However, the removal of the W atoms from the surface costs anywhere between 1.00 and 4.42 eV. We also find that the removal of WOCl₄ as a non-W_s etch product costs less energy when compared to WO₂Cl₂ since the former removed only one surface-bound O atom whereas the latter removes two, which costs more energy. This is further supported by the observation that, whenever a WOCl₄ species is formed on the surface, it readily desorbs from the surface. In contrast, WO₂Cl₂ is bound to the surface primarily as a result of the strong W–W and W–O bonds. Therefore, WO₂Cl₂ should become volatile only at elevated temperatures.

To consider other combinations of etch products, 2 WCl₆ adsorption on the O-covered W surface is studied as shown in Figure 6. Similar to the low coverage case, the interaction

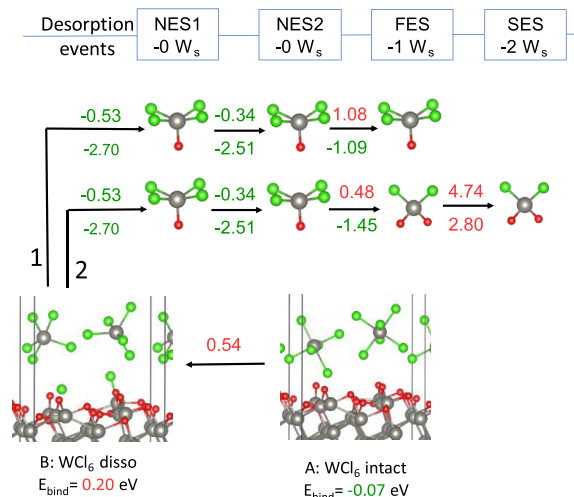


Figure 6. Minimum energy geometries showing the non-volatile adsorption of 2 WCl₆ on the O-covered W surface and energies needed to remove volatile species from them. The values in red indicate energy loss, and the values in green indicate energy gain. The values above the arrow are the desorption energies of the respective volatile species, while the values below are the free energy change of desorption at 500 K.

between WCl₆ and the substrate was weak with a binding energy of just −0.07 eV/WCl₆ (geometry A). Dissociation of the adsorbed WCl₆ molecules by donating one Cl atom to the surface (geometry B) costs about 0.27 eV/WCl₆. Keeping geometry B as the starting point, we computed the desorption energies for the removal of three consecutive WOCl₄ species. The removal of the first two volatile species (NESs) is exoergic, while the removal of the third WOCl₄ (FES) has a cost of 1.08 eV. The removal of the above three volatile products was favorable at 500 K. In another desorption sequence, we studied the desorption of two consecutive WOCl₄ as the NES followed by two WO₂Cl₂ as the FES and SES. While the removal of the first two WOCl₄ species was essentially the same as in the previous case, the subsequent removal of the first WO₂Cl₂ species costs less energy compared to the energy cost for removal of WOCl₄ in the other sequence. However, the removal of the second surface W atom in the form of WO₂Cl₂ has a huge energy penalty of 4.7 eV. We find that the removal of a surface atom as WOCl₄ is not more energetically favorable than its removal as WO₂Cl₂, which is understandable as more surface W–O/W–W bonds have to be broken to form the former than the latter in this particular

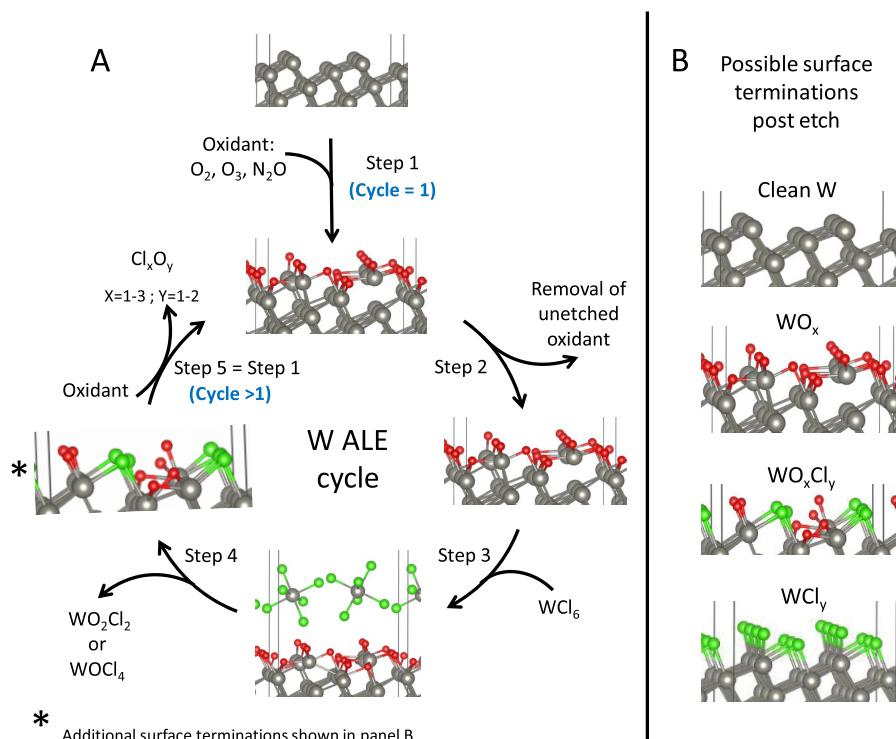


Figure 7. Panel (A) shows a schematic representation of an idealized W ALE cycle. The energies of the individual steps are listed in Table 2. Atom color coding is as follows: W = gray, O = red, and Cl = green. Step 1 depicts the initial functionalization of the surface with the desired oxidant. Step 2 describes the purging of the excess oxidant from the reaction chamber. Step 3 describes the introduction of WCl_6 , which exhibits no spontaneous dissociation until elevated temperatures. Step 4 depicts the liberation of volatile WO_xCl_y species as the most likely etch species. Step 5 in an idealized system results in the repassivation of the surface leading into the next cycle and is thus labeled as equivalent to step 1. Additional passivated surface post etch is shown in panel (B). WCl_x terminated surface is discussed in detail in Figure 8.

case. On the other hand, it is more favorable to form $WOCl_4$ as a NES from the incoming WCl_6 molecule when compared to WO_2Cl_2 .

Now, let us look at the chlorination pulse when Cl_2 is used as the gas phase reactant. Similar to WCl_6 , Cl_2 did not adsorb strongly via dissociation at the O-rich region of the WO_x surface. A binding energy of just -0.05 eV per Cl_2 molecule is obtained. A short MD simulation at 800 K also failed to capture Cl_2 dissociation as the Cl_2 molecule simply diffused away from the surface due to the repulsion between the Cl and surface O atoms. Therefore, we studied the dissociative adsorption of 1 and 2 Cl_2 on the model WO_x surface, which released an energy of -3.11 eV to -3.92 eV per adsorbed Cl_2 . The corresponding geometries are schematically presented in section S8 of the Supporting Information. From the above structures, desorption energies of volatile species such as WO_2Cl_2 and $WOCl_4$ are computed to be 1.33 and 0.73 eV, respectively. It is to be noted that there is no NES when Cl_2 is used in the chlorination pulse. Also, as mentioned earlier in Section 3.1, each Cl_2 molecule is capable of forming, at most, one etch product (WO_2Cl_2 as the FES), whereas two Cl_2 molecules are needed to form one $WOCl_4$ (as the FES) or two WO_2Cl_2 (as the FES and SES). On the bare W surface, Cl_2 dissociates spontaneously at low temperatures (discussed in detail in section S4 of the Supporting Information). Homolytic dissociation of Cl_2 is required on the WO_x surface so that volatile species can be formed. Dissociation of Cl_2 in vacuum requires 3.1 eV, while the dissociation of one W–Cl bond from the WCl_6 molecule in vacuum required a slightly lower 2.5 eV. Based on the above, WCl_6 may dissociate at lower temper-

atures when compared to Cl_2 . Therefore, in the next section, we will discuss the full ALE cycle with WCl_6 as the second pulse chemical.

3.4. The Full Atomic Layer Etch Cycle. In this section, we discuss the full cycle of the W ALE process, with O_2 , O_3 , and N_2O as the sfirst pulse chemical and WCl_6 as the second pulse chemical, which is summarized in Figure 7A. Steps 1, 2, 3, and 4 correspond to pulse-1, purge-1, pulse-2, and purge-2 of the W ALE cycle, respectively. Step 5 is also the step 1 of the second and subsequent ALE cycles, whereas steps 2, 3, and 4 are common for all ALE cycles. The first ALE pulse is the adsorption of the oxidizing agents (O_2 , O_3 , and N_2O). In the oxidation step, step 1 of the first ALE cycle, energy is gained by the dissociative adsorption of the reactant molecules on the bare W surface. We will show later that subsequent ALE cycles do not start with a bare W surface but rather a WO_xCl_y surface. We show in Table 2 and section S3 of the Supporting Information that this step realizes a significant enthalpy gain by the formation of stable W–O bonds, but the associated entropic loss was comparatively too small to make the reaction unfavorable at ALE relevant temperatures.

An energy of -4.31 to -4.80 eV (-3.61 to -4.63 eV at 500 K), depending on the first pulse reactant used, is gained per surface W atom in the supercell (W_s , note that there are 12 W_s in the supercell) in this step. This amounts to an energy of -3.69 to -4.11 eV released per O atom adsorbed. In step 2, unreacted reactant molecules are purged from the etch chamber along with any byproducts formed.

In step 3, the dissociative adsorption of WCl_6 is either mildly favorable energetically (-0.08 eV/ W_s) or costs energy (0.12

Table 2. Energy Requirements (Ranges with Most Favorable (MF) and Least Favorable (LF) Values in eV) for the Individual Steps of the ALE Cycle in Figure 7^a

step	ΔE (eV)		ΔG (eV) 500 K		per
	MF	LF	MF	LF	
O ₂ /O ₃ /N ₂ O pulse	−4.80	−4.31	−4.63	−3.61	W _s
purge 1	0.00	0.00	0.00	0.00	W _s
WCl ₆ pulse	−0.08	0.12	0.11	0.30	W _s
	−4.88	−4.19	−4.52	−3.31	W _s
purge 2	−0.53	0.57	−2.70	−1.37	NES
	0.48	2.72	−1.45	0.64	FES
	4.42	4.74	2.48	2.80	SES

^aHere, W_s refers to a surface W atom. NES, FES, and SES are acronyms for non-W_s etch, first-W_s etch, and second-W_s etch species, respectively. The values in bold refer to the cumulative values combining oxidation pulse and chlorination pulse. For the free energy values at 500 K, the contributions from surface models are not included as they were estimated to be less than 4% of the total value.

eV/W_s). Moreover, this step is endergonic at 500 K mainly due to the entropic loss. However, the cumulative energetic gain through steps 1 to 3 is still exoergic, −4.19 to −4.88 eV per W_s (−3.31 to −4.52 eV per W_s at 500 K). We use the energy per W_s since it is convenient to compare to the energy required to desorb the etch products in the next step. As discussed earlier, we have also performed a short 2.5 ps MD study and observed spontaneous dissociation of the WCl₆ molecule on the WO_x surface at 800 K. We chose 800 K to accelerate the MD simulation. However, this process is also possible from simulations at temperatures close to the ALE relevant temperature, e.g., 500 K, provided long trajectories are computed, which are computationally expensive.

In step 4, the formation and desorption of the non-W_s etch and W_s etch byproducts (*X* number of WO₂Cl₂ and *Y* number of WOCl₄) take place. The first product per incident WCl₆ is a NES, and depending on the stoichiometry of the byproduct, there could additionally be one (FES) or two (FES and SES) W_s etch species. We have determined the typical energy requirement for the removal of the volatile NES, FES, and SES as shown in Table 2. The NES required the least energy (sometimes this is exoergic) as evident from the negative free energies at 500 K. The removal of the FES requires between 0.48 and 2.72 eV, and it can even be spontaneous at higher temperatures (0.64 to −1.45 eV at 500 K) due to net entropic gain. However, the removal of a SES requires between 4.42 and 4.47 eV of energy and it is also not favorable at high temperatures. However, the cumulative enthalpic gain per surface W atom up to this point might compensate these energy requirements for the removal of the FES and the SES at elevated temperatures. Thus, the removal of the NES is spontaneous while the removal of the FES and the SES costs energy and is subject to free energy changes in the cycle, the latter more so than the former. Since we gain a free energy of −39.72 to −54.24 eV per supercell (12 W_s) at 500 K from the first three steps, the overall energy change for the entire cycle including the removal of NES, FES, and SES from a WCl₆ molecule will still be favorable (negative free energy).

At this point, the second and subsequent ALE cycles will start from a surface that is potentially one of the four shown in Figure 7B, that is WO_xCl_y, WCl_y, or WO_x terminated or a potentially bare W surface. The reason for this is that the surface geometry after the WCl₆ pulse will be dependent on

the process conditions and the reactions that take place in the second pulse and in the second purge steps. The W surface is not likely to be terminated with bare tungsten as this is only possible when all the adsorbed WCl₆ species form volatile WO_xCl_y species, thereby depleting the adsorbed Cl and all surface O atoms. WO_x will only result if the WCl₆ molecules in the second pulse do not dissociatively react with the substrate and cause an etch delay, which is also not very probable. It is more likely that the surface will now be partially covered with Cl species (WCl_x) because the WCl₆ and Cl₂ molecules dissociate readily on the bare W surface, which is now exposed after volatile products are removed from the WO_x surface, by releasing −4.5 and −5.9 eV, respectively. It is also possible that there are some O atoms remaining on the surface forming non-volatile WO_xCl_y species as suggested by Xie *et al.*^{25,26} Our calculations also predict that the removal of the SES requires a very high energy, which indicate a O_xCl_y termination of the W surface. Therefore, as subsequent ALE cycles proceed, the oxidizing agents will have to react with an (oxy)chloride termination of the W surface rather than a bare W surface and this may influence the overall etch rate.^{25,26}

To understand the potential impact of this, we studied a model system of six O₂ molecules per supercell interacting with a fully Cl-covered W surface that may be present after one etch cycle, as shown in Figure 8. When initially intact O₂ molecules are adsorbed on the Cl-covered W surface (Figure 8a), no spontaneous reaction is observed due to the repulsion between the Cl and O atoms. We note that O₂ dissociates spontaneously on Cl-free regions on the W surface, though we cannot see that in our model because of the full surface passivation with Cl. We therefore examined two configurations

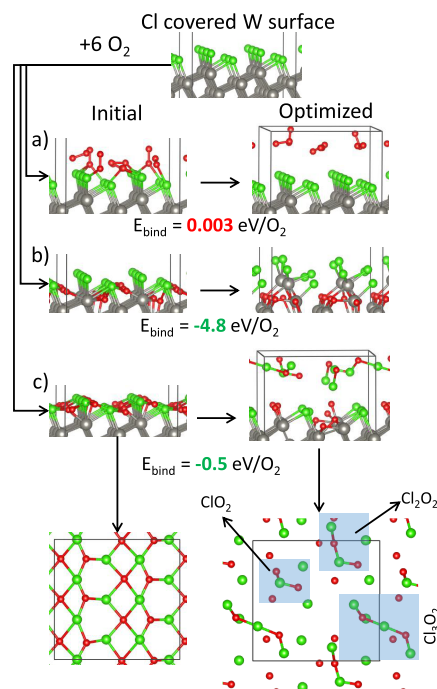


Figure 8. Interaction of O₂ on the Cl-covered W surface, which is identified as one of the possible surface terminations at the end of the ALE cycle in Figure 7. In case (a), six intact O₂ molecules are adsorbed on the surface with 0 eV/O₂. In (b) and (c), 12 O atoms (six dissociated O₂ molecules) are adsorbed close to surface W and close to Cl atoms with a release of −4.8 and −0.5 eV/O₂, respectively. The color coding is as follows: W = gray, O = red, and Cl = green.

whereby dissociated O_2 molecules are adsorbed on the Cl-covered W surface (Figure 8b,c). In Figure 8b, the resulting O atoms are adsorbed close to the surface W atoms, which on relaxation resulted in the spontaneous formation of surface-bound WO_2Cl_2 species with a release of -2.4 eV per adsorbed O atom. In the second configuration (Figure 8c), oxygen atoms are adsorbed close to the Cl atoms, and upon relaxation, we observe the spontaneous formation and desorption of oxychloride species such as ClO_2 , Cl_2O_2 , and Cl_3O_2 , with a gain of -1.0 eV/O. This suggests that the incoming O_2 molecule in the second ALE cycle must first dissociate before adsorbing at the fully chlorinated surface in order to react with it. As mentioned earlier, this pre-dissociation of the O_2 molecules is not needed when the Cl ligands are sparsely distributed on the surface.

4. CONCLUSIONS

A first-principles investigation of a thermal atomic layer etch process for metallic tungsten using sequential oxidation and chlorination is presented. O_2 , O_3 , and N_2O were examined as the oxidizing agents to modify the oxidation state of the surface W atoms. WCl_6 and Cl_2 were examined as ligand exchange agents for the material removal pulse. The reactant molecules chosen for both ALE pulses meet the requirement of being stable in the gas phase at the ALE operating temperature. Runaway oxidation of W is possible when O_2 and O_3 are used in the first ALE pulse. However, a thin layer of WO_x can be obtained by carefully controlling the reactant pressure, temperature, and exposure time. The most important aspect of the ALE process is the self-limiting nature of the reactions taking place in the two pulses. Qualitative predictions based on thermochemical analysis showed that bulk oxidation was hindered by thermodynamic barriers at ALE relevant temperatures and pressures. Similarly, spontaneous etching of W by WCl_6 and Cl_2 was predicted to be unfavorable, even though these molecules spontaneously react with the bare W surface.

The mechanisms of the dissociative adsorption of the oxidizing agents on the bare W surface were reported. Exposure to O_2 , O_3 , and N_2O resulted in an O-covered surface with ozone being the most reactive chemical for the first pulse followed by O_2 and N_2O . To keep the computational effort tractable, we considered a monolayer oxidized surface model of W at the end of the first half ALE cycle. However, WCl_6 molecules interacted weakly with the oxidized W surface due to net repulsion between the O and Cl atoms and also because of the kinetic barriers inhibiting the dissociation of WCl_6 molecule. Thermal energy is required to dissociate the WCl_6 molecule so that it can donate its Cl ligands to the surface leading to the formation of the target volatile species, primarily $WOCl_4$ and WO_2Cl_2 . The incoming WCl_6 molecule can form $WOCl_4$ as a non- W_s etch species by exchanging two Cl ligands for one O atom from the surface. This also results in the formation of a WO_2Cl_2 species on the surface. The desorption of WO_2Cl_2 has a relatively higher energy cost as compared to $WOCl_4$ due to the relatively high coordination of the W atom in surface-bound WO_2Cl_2 . The other possibility for the incoming WCl_6 molecule is to dissociate completely by donating all of the Cl ligands to the surface and form three WO_2Cl_2 molecules as volatile species. Therefore, each WCl_6 molecule is capable of removing up to two surface W atoms with a maximum energy cost of 4.42 eV. The desorption of the first- W_s etch product costs between 0.48 and 2.72 eV, which is thermodynamically favorable in the

overall energy cycle, whereas the desorption of the second etch product costs almost 5 eV and is unfavorable.

Due to the favorable reaction of WCl_6 and Cl_2 on the bare W surface, it will be covered with O and Cl ligands at the end of the first ALE pulse. In the second ALE cycle, the oxidizing agents must remove the Cl ligands first, either by forming volatile W oxychlorides or stable oxychloride chains, before accessing the surface W atoms, which might slow down the etch process. Computational investigations of such ALE processes provide the required understanding of the mechanism of etch reactions. Such understanding is vital to the design of new and robust ALE processes in the future.

■ ASSOCIATED CONTENT

Supporting Information

The Supporting Information is available free of charge at <https://pubs.acs.org/doi/10.1021/acsami.0c06628>.

Atomic models of W bulk and surface, schematics of second ALE pulse, mechanism of N_2O pulse, Natarajan–Elliott “N-E” analysis of the oxidation pulse, reaction of WCl_6 and Cl_2 on the bare W surface, dissociation pathway of WCl_6 , description of thermal ALE of metals, influence of reactant and product pressures, and Cl_2 in material removal pulse (PDF)

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Notes

The authors declare no competing financial interest.

■ ACKNOWLEDGMENTS

Authors S.K.N. and M.N. thank Intel Corporation for funding. They also thank the Irish Centre for High-End Computing (project code: tiche081c) and Science Foundation Ireland-funded computing cluster at Tyndall for the computer time.

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